# Dynamics of a Crash Victim—A Finite Segment Model

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A vehicle-occupant, crash-simulation model is developed. The model consists of a finite-segment representation of the human body, together with a vehicle cockpit that includes a seat and crash intrusion surfaces (windshield, doors, etc.). The human body model is constrained to the cockpit by seat and shoulder belts. The equations of motion are developed for the model by using Lagrange's form of d'Alembert's principle. These equations then are coded into a computer program and solved numerically using a fourth-order Runge-Kutta method. Sample motions for several crash situations are presented and discussed.

#### Nomenclature

= acceleration of  $G_k$  in R $B_k$   $F_k$   $F_p$   $G_k$   $I_k$ = bodies of the human model (k=2,...,12)= external force applied to  $B_k$ = generalized active force = generalized inertia force = mass center of  $B_k$ =inertia dyadic of  $B_k$  relative to  $G_k$  $m_k$ = mass of  $B_k$  $M_k$ = external torque applied to  $B_k$  $O_k$ = reference point of  $B_k$  $R^{k}$ = mass center  $(G_k)$  position vector in  $B_k$ =inertial reference frame = partial rate of change of position  $v_{kji}$ = generalized coordinates (i = 1, ..., 48)= angular acceleration of  $B_k$  in R= reference point  $(O_k)$  position vector = angular velocity of  $B_k$  in R

#### Introduction

DURING the past decade and especially during the past five years, there has been considerable research effort expended upon the development of finite-segment and finite-element models of structures and various dynamical systems. Of all of these models, perhaps the most extensive and imaginative effort has been in the development of finite-segment models of the human body and its dynamics, particularly in crash or high-acceleration environments. References 1-21 provide a summary of this research and its progress to date, with Ref. 1 providing a detailed comparison of the efforts. It suggests that the most elaborate models are those of Ref. 15 and 18. Success in these efforts has been related closely (although not exclusively) to advances in digital computer technology.

In this paper, the theoretical elements and basic features of a modern finite-segment, human-body, crash-victim model are presented. The paper is divided into five parts, with the first part describing the modeling process itself and the various associated theoretical problems. The second part presents the basic kinematics and dynamics of the model. The

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Index categories: Aircraft Cabin Environment and Life Support Systems; Computer Technology and Computer Simulation Techniques; Structural Dynamics Analysis. third part briefly discusses the specific features of the UCIN code, with some specific examples and illustrations presented in the fourth part. The final part of the paper presents some conclusions regarding the range of application of the model and some ideas regarding future model development.

## **Finite-Segment Modeling**

Finite-segment modeling in general, and more specifically human-body modeling, is perhaps most efficiently discussed in the context of modeling and of "general chain systems." A general chain system is defined to be a series of rigid bodies connected to each other in treelike fashion such that adjoining bodies share a common point and such that no closed loops are formed. Figure 1 shows an example of such a system.

The modeling and dynamics of general chain systems have been studied by Huston and Passerello, <sup>22,23</sup> Roberson, <sup>24</sup> Likins, <sup>25-27</sup> Chace and Bayazitoglu, <sup>28</sup> Hooker, <sup>29</sup> Fleischer, <sup>30</sup> and others. The basic difficulty encountered in all of these studies is the complexity of the geometry. This, in turn, presents difficulties in finding a kinematical description of the system and in obtaining the governing dynamical equation of motion. Huston and Passerello<sup>22</sup> have shown that the kinematics can be described conveniently and efficiently by using local coordinates (that is, angles between adjoining bodies) as opposed to absolute coordinates (that is, orientation angles in space) and by using a combined vectormatrix-tensor approach. They also maintain that using

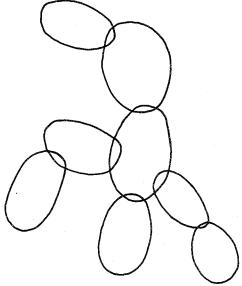


Fig. 1 A general chain system.

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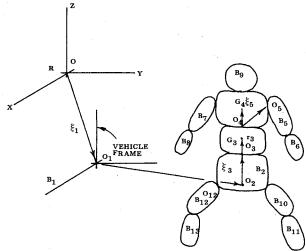


Fig. 2 The human-body model with a vehicle frame.

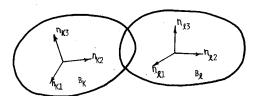


Fig. 3 Two typical adjoining bodies of the model.

Lagrange's form of d'Alembert's principle<sup>31-34</sup> is the most efficient way to obtain the governing dynamical equations of motion. They point out that Lagrange's form of d'Alembert's principle provides the exact number of governing equations (like Lagrange's equations, but not Newton's laws) and that it allows one to avoid lengthy involved differentiation of energy functions (like Newton's laws, but not Lagrange's equations). That is, for large complex systems, Lagrange's form of d'Alembert's principle incorporates the relative advantages of Lagrange's equations and Newton's laws while avoiding the corresponding disadvantages.

In this paper, the methods of Huston and Passerello are used to develop the governing equations of motion for a chain system representing a human-body model in a cockpit or vehicle frame as shown in Fig. 2. In this figure, the 12 rigid bodies are a spheriod, elliptical cylinders, and frustrumes of elliptical cones representing the various limbs of the human body. The system contains 42 degrees of freedom (6 for the translation and rotation of the vehicle frame, 6 for the translation and rotation of a reference body of the model relative to the vehicle frame, and 33 for the orientation of the 11 remaining bodies relative to their adjacent lower numbered bodies). The following section outlines the procedures used in obtaining the equations of motion for this system.

## **Kinematics and Equations of Motion**

In Fig. 2, R represents an inertial reference frame, B represents the vehicle frame, and  $B_2...,B_{I3}$  represents the various limbs of the human-body model as shown.  $B_2$  is the reference body of the model. O and  $O_1$  are the origins (reference points) of R and  $B_1$ .  $O_2$  is an arbitrarily chosen reference point of  $B_2$ . The reference points of each of the remaining bodies are the connecting points with the adjacent lower numbered bodies (e.g.,  $O_3$ ,  $O_4$ ,  $O_5$ ,..., $O_{10}$ ). The vectors  $\xi_k$  (k=1,...,13) locate the reference points of the adjacent lower numbered bodies. (Note that, except for  $\xi_1$  and  $\xi_2$ ,  $\xi_k$  is fixed in the adjacent lower numbered body.) The vectors  $r_k$  (k=2,...,13) locate the mass centers  $G_k$  of  $B_k$  relative to  $O_k$ .

Consider a typical pair of adjoining bodies of the model such as  $B_k$  and  $B_s$  as shown in Fig. 3. Then the orientation of-

 $B_i$  relative to  $B_k$  may be defined in terms of the relative inclination of the dextral, orthogonal unit vector sets  $n_{ki}$  and  $n_{ii}$  (i=1,2,3), which are fixed, respectively, in  $B_k$  and  $B_i$ . This is done as follows. Let  $B_k$  and  $B_i$  be oriented so that  $n_{ki}$  and  $n_{ii}$  are respectively parallel. Then  $B_i$  may be brought into any given orientation relative to  $B_k$  by three successive dextral rotations about axes parallel to  $n_{ki}$ ,  $n_{k2}$ , and  $n_{k3}$  through the orientation angles  $\alpha_{ki}$ ,  $B_{ki}$ , and  $\gamma_{ki}$ ,  $n_{ki}$  and  $n_{ii}$  then are related to each other by the expressions§

$$n_{ki} = SKL_{im}n_{\ell m} \tag{1}$$

where SKL is a  $3 \times 3$  orthogonal tranformation matrix defined

$$SKL_{im} = n_{ki} \cdot n_{im} \tag{2}$$

SKL may be written as the product of three orthogonal matrices as

$$SKL = \alpha KL \ \beta KL \ \gamma KL \tag{3}$$

where

$$\alpha KL = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c\alpha_{k\ell} & -s\alpha_{k\ell} \\ 0 & s\alpha_{k\ell} & c\alpha_{k\ell} \end{bmatrix}$$
(4a)

$$\beta KL = \begin{bmatrix} c\beta_{k\ell} & 0 & s\beta_{k\ell} \\ 0 & 1 & 0 \\ -s\beta_{k\ell} & 0 & c\beta_{k\ell} \end{bmatrix}$$
 (4b)

$$\gamma KL = \begin{bmatrix} c\gamma_{k\ell} & -s\gamma_{k\ell} & 0 \\ s\gamma_{k\ell} & c\gamma_{k\ell} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$
 (4c)

where  $s\alpha_{kl}$  and  $c\alpha_{kl}$  represent the sine and cosine of  $\alpha_{kl}$ . From Eq. (1), it easily is seen that, with three bodies  $B_k$ ,  $B_l$ , and  $B_m$ , the shifter transformation matrix obeys the following chain and identity rules:

$$SKM = SKL SLM \tag{5}$$

and.

$$SKK = I = SKL SLK = SKL SKL^{-1}$$
 (6)

These shifter matrices are used to transform the components of vectors and tensors referred to one body of the system into components referred to any other body of the system and, in particular, to the inertial frame R. It is necessary then also to have an expression for the derivative of the shifter matrix, especially the matrix SOK, where the O refers to the inertial frame R. Huston and Passerello<sup>22,23</sup> have shown that the derivative may be calculated by a simple matrix product as

$$SOK = WOK SOK$$
 (7)

where WOK is defined as

$$WOK_{im} = -e_{imn}\omega_n^k \tag{8}$$

<sup>§</sup>Regarding notation, repeated subscripts such as m in the right side of Eq. (1) represent a sum over the range (1,...,3) of that index.

where the  $e_{imn}$  is the standard permutation symbol<sup>35</sup> and  $\omega_n^k$ are the components of the angular velocity of  $B_k$  referred to unit vectors  $n_{on}$  fixed in R.

The angular velocity  $\omega^k$  of  $B_k$  in R is obtained from the addition formula34

$$\omega^k = \omega^I + {}^I\omega^2 + \dots + {}^J\omega^k \tag{9}$$

where  ${}^{j}\omega^{k}$  is the angular velocity of  $B_{k}$  relative to  $B_{j}$ . The series in Eq. (9) proceeds outward from  $B_1$  to  $B_k$  through the branch of the chain or model containing  $B_k$ . From Fig. 2 and the preceding analysis, it easily is seen that  $j\omega^k$  may be expressed as

$$^{j}\omega^{k} = SOJ_{im} (\dot{\alpha}_{jk}\delta_{ml} + \beta_{jk}\alpha JK_{m2} + \gamma_{jk}\alpha JK_{ml}\beta JK_{\ell 3}) \mathbf{n}_{oi}$$

$$(10)$$

where  $\delta_{mn}$  is Kronecker's delta symbol or identity tensor, <sup>35</sup> and  $n_{oi}$  (i = 1,2,3) are unit vectors fixed in R. By repeatedly substituting Eq. (10) into Eq. (9),  $\omega^k$  can be seen to have the

$$\boldsymbol{\omega}^k = \omega_{kpm} \dot{\boldsymbol{x}}_p \boldsymbol{n}_{om} \tag{11}$$

where the  $\dot{x}_p$  (p = 1,...,42) are the generalized coordinates representing the 42 degrees of freedom. By using Eq. (10), it is seen that the nonzero terms of  $\omega_{kpm}$  take on of the following forms:

$$\omega_{kpm} = \begin{array}{c} SOJ_{ml} \\ \omega_{kpm} = SOJ_{mn} \alpha JK_{n2} \\ SOJ_{mn} \alpha JK_{nh} \beta JK_{h3} \end{array}$$
(12)

depending upon whether  $x_p$  is the first, second, or third dextral angle  $(\alpha, \beta, \text{ or } \gamma)$  defining the orientation of  $B_k$  relative to its adjacent lower numbered body. Also, it is shown easily that, for any two bodies  $B_k$  and  $B_l$  in the same branch of the model,  $\omega_{\ell_{pm}} = \omega_{kpm}$  for  $\ell > k$  if  $\omega_{kpm} \neq 0$ . The angular acceleration  $\alpha^h$  of  $B_k$  in R may be obtained by

differentiating in Eq. (11), that is,

$$\alpha^{lsk} = (\omega_{kpm} \ddot{x}_p + \dot{\omega}_{kpm} \dot{x}_p) n_{om}$$
 (13)

where  $\dot{\omega}_{kpm}$  may be obtained directly by differentiating in Eq.

The velocity  $v^k$  of  $G_k$ , the mass center of  $B_k$ , relative to Rmay be obtained by differentiating the position vector  $p_k$ locating  $G_k$  relative to O. From Fig. 2,  $p_k$  is seen to have the

$$p_i = \xi_1 + \xi_2 + (\Sigma SOJ_{mn}\xi_n^j + SOK_{mn}r^k)n_{om}$$
 (14)

where  $\Sigma$  indicates a sum over J (or j) for the bodies in the branch containing  $B_k$  (j < k). Performing the differentiation in Eq. (14) shows that  $v^k$  may be written in a form analogous to Eq. (11) as

$$\boldsymbol{v}^k = \boldsymbol{v}_{kpm} \dot{\boldsymbol{x}}_p \boldsymbol{n}_{om} \tag{15}$$

where the  $v_{kpm}$  are obtained through using Eqs. (7) and (8).

The acceleration  $a^k$  of  $G^k$  in R may be obtained by differentiating in Eq. (15). That is,

$$\boldsymbol{a}^{k} = (\boldsymbol{v}_{kpm} \ddot{\boldsymbol{x}}_{p} + \dot{\boldsymbol{v}}_{kpm} \dot{\boldsymbol{x}}_{p}) \boldsymbol{n}_{om} \tag{16}$$

The governing dynamic equations of motion now may be obtained as follows. Consider again the model shown in Fig. 2. Let the externally applied forces be replaced by a system of forces consisting of single forces  $F_k$  passing through  $G_k$  (k=1,2,...,13) together with couples with torques  $M_k$  applied to  $B_k$ . Then Lagrange's form of d'Alembert's principle states that the governing equations of motion of the system are 34

$$F_p + F_p^* = 0 \ (p = 1, ..., 42)$$
 (17)

where  $F_n$  is called the generalized active force and is given by

$$F_p = v_{kpm} F_{km} + \omega_{kpm} M_{km} \tag{18}$$

where  $F_{km}$  and  $M_{km}$  are the  $n_{om}$  components of  $F_k$  and  $M_k$  (k=2,...,13).  $F_p^*$  is called the generalized inertia force and is

$$F_p^* = v_{kpm} F_{km}^* + \omega_{kpm} M_{km}^*$$
 (19)

where  $F_{km}^*$  and  $M_{km}^*$  are n are  $n_{om}$  components of the inertia forces  $F_k^*$  and inertia torques  $M_k^*$ , which, in turn, are  $^{34}$ 

$$F_k^* = -m_k a^k \qquad \text{(no sum)} \tag{20}$$

and

$$M_k^* = -I_k \cdot a^k - \omega^k x (I_k \cdot \omega^k)$$
 (no sum) (21)

where  $m_k$  is the mass of  $B_k$ , and  $I_k$  is the inertia dyadic of  $B_k$ relative to  $G_k$ . Hence, by substituting from Eqs. (9-21), the equations of motion may be written in the form

$$a_{pq}\ddot{x}_q = f_p \ (p, q = 1, ..., 42)$$
 (22)

where  $a_{pq}$  and  $f_p$  are

$$a_{pq} = m_k v_{kpm} v_{kqm} + I_{kmn} \omega_{kpm} \omega_{kqn}$$
 (23)

$$f_{p} = -(F_{p} + m_{k}v_{kpm}v_{kqn}\dot{x}_{q} + I_{kmn}\omega_{kpm}\dot{\omega}_{kqn}\dot{x}_{q}$$

$$= e_{hmn}\omega_{kph}\omega_{kqm}\omega_{kri}I_{kni}\dot{x}_{q}\dot{x}_{r})$$
(24)

where  $I_{kmn}$  are the  $n_{om}$  components of  $I_k$ .

Equations (22) form a set of 42 simultaneous ordinary, nonlinear differential equations determining the 42 generalized coordinates  $x_p$  of the system. If some (or all) of the  $x_p$  are specified, then the differential equations become algebraic equations for the unknown forces or moments associated with the specified  $x_p$ . Since the coefficients  $a_{pq}$  and  $f_p$  of these equations are algebraic functions of the physical parameters and the generalized coordinates, the equations are ideally suited for generation on a digital computer.

#### **UCIN Code**

Algorithms have been written to develop the governing equations with a computer. These algorithms form the basic structure of the UCIN vehicle-occupant, crash-simulation code.21 A number of special features then are added to simulate a vehicle-occupant system. These include 1) a vehicle (or cockpit) that can undergo up to six (three translation and three rotation) simultaneous accelerations; 2) an occupant seat rigidly attached to the vehicle frame consisting of up to seven springs and dampers of arbitrary moduli supporting the head, the upper middle and lower back, the tail, and the thighs; 3) up to 10 seat belts arbitrarily attached to the vehicle frame and the bodies of the human model; these belts are modeled as linear springs with arbitrary moduli; 4) arbitrary angle restraints at each joint of the human body model; these consist of dampers with arbitrary moduli; and 5) intrusion surfaces consisting of 13 planes of arbitrary position and orientation modeling the cockpit and providing a basis for "hits" of the human-body model with the cockpit. The governing equations then are solved numerically using a fourthorder, variable-step, Runge-Kutta integration routine. Input for the computer program consists of the physical data for the human-body model (masses and inertia matrices of the bodies, mass center position, and joining locations), the moduli for the various springs and dampers (see the foregoing), the number and attach points of the various seat belts, the cockpit intrusion surface geometry, and the vehicle

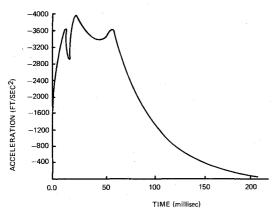


Fig. 4 Sled deceleration.

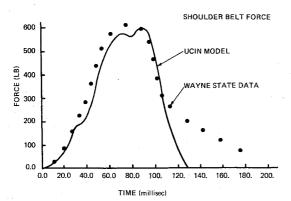


Fig. 5 Comparison of shoulder belt forces.

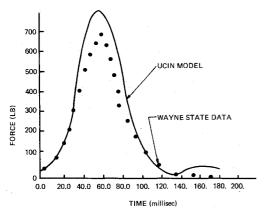


Fig. 6 Comparison of vertical lap belt forces.

acceleration profile. The output consists of position, velocity, and acceleration profiles of the mass centers and joints relative to both the vehicle frame and inertia space.

The following section contains some illustrations of the use of the code. In each of these, the physical data input for the human-body model was taken from Hanavan<sup>36</sup> for his 50 percentile Air Force mean man. Either the other input data (springs and dampers) were selected to match those used in the experimental configuration, or moderate (i.e., average) values were assumed when no other source was available.

## **Verification and Examples**

There is little experimental data available to date which can be used to check or verify the preceding computer code and others like it. However, an attempt at verification was made for some data gathered by King et al.<sup>37</sup> In these experiments, they placed a cadaver in a sled that impacted a stop, generating a sled deceleration profile, as shown in Fig. 4. The shoulder belt, vertical lap belt, and seat pan forces were

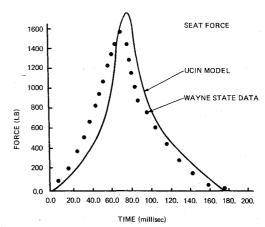


Fig. 7 Comparison of seat forces.

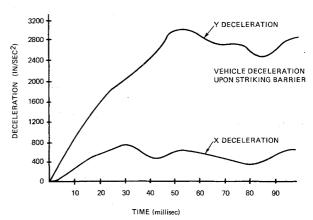


Fig. 8 Vehicle deceleration.

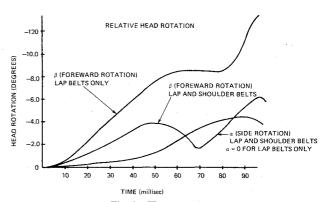


Fig. 9 Head rotation.

measured. These forces then were calculated using the UCIN code with approximately the same sled (vehicle) deceleration profile and lap and shoulder belt configuration. (The acceleration profile was approximated by 25 straight line segments.) A comparison of the results is shown in Figs. 5-7.

In a second example, experimental data from a vehicle striking a guard rail or roadside barrier were used as input for the UCIN code. The specific acceleration profile and the resultant displacement of the head and chest using both lap belts and a combination of lap and shoulder belts is shown in Figs. 8-10.

Finally, in an attempt to measure the relative effectiveness of lap and shoulder belts, a run was made simulating a frontend collision of a vehicle. The head pitch of the model (forward rotation relative to the chest) was calculated using both lap belts and a combination of both lap and shoulder belts, as normally used in automobiles. The results, shown in Fig. 11,

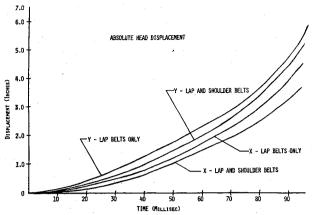


Fig. 10 Head displacement.

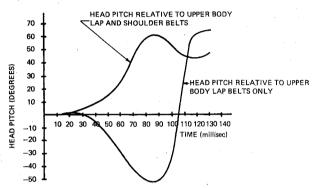


Fig. 11 Comparison of lap and shoulder belts.

clearly illustrate a "whiplash" effect when only lap belts are used.

## **Discussion and Conclusions**

The analysis shows that it is now possible to construct reliable, finite-segment simulation models of crash vehicle occupants. The success of the modeling is due primarily to recent advances in digital computing machines and to new approaches in formulating the governing dynamical equations of motion.

The examples presented show the range of applicability of the model. In the particular cases studied, the advantages of combined shoulder and lap belt restraints as opposed to lap belts alone are remarkable, particularly in the prevention of "whiplash."

Finally, it appears that future work will include refinements of the model, particularly in vehicle (cockpit) and restraint modeling, model verification, and in making provision for variable numbers of human-body segments. Application of the model in the design of vehicle interior and safety devices is clear but is yet to be developed.

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